

Gamma rays from cosmologically distant sources of UHECR

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Outline

- Cosmologically distant sources
- Gamma-ray traces of UHECR and IGMF
- Synchrotron gamma rays as signature of UHECR
- Gamma rays from distant blazars as signature of CR
- Conclusion

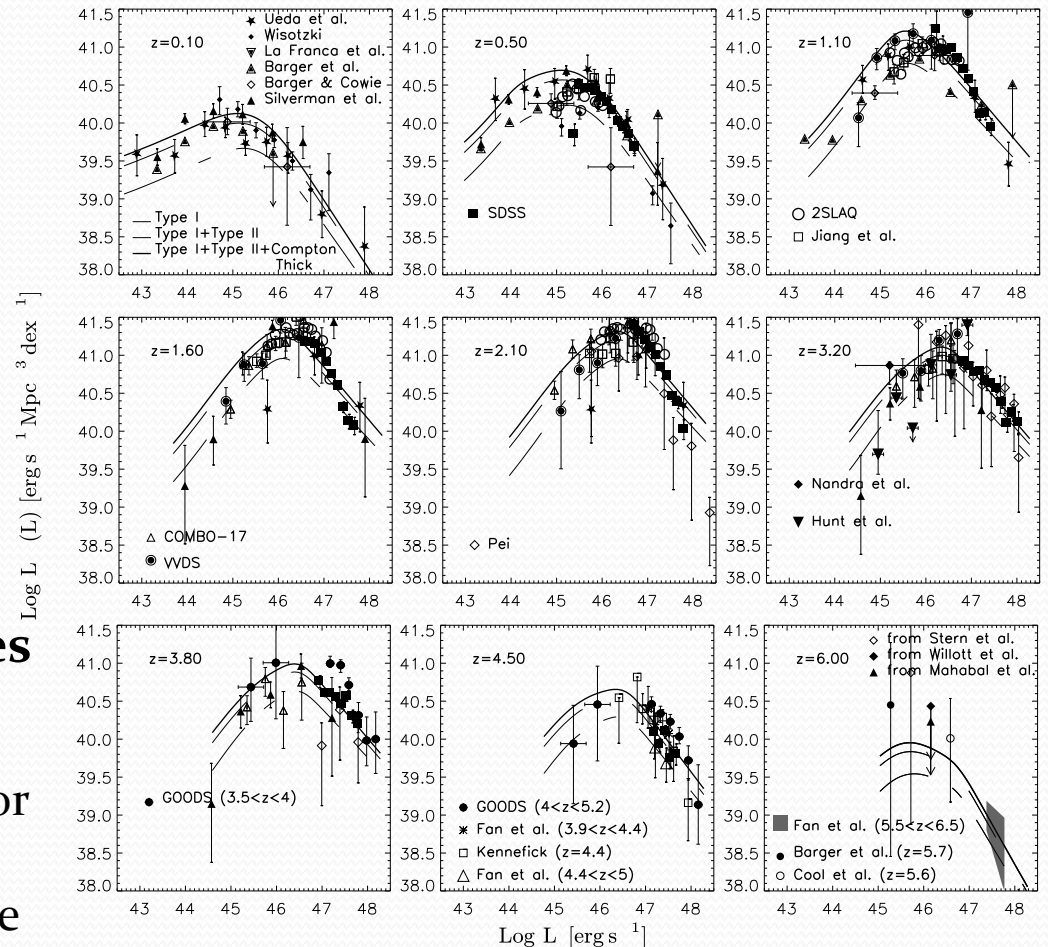
Why cosmologically distant sources?

With increase of redshift:

- Energetic increases
- Probably more effective accelerators
- Number of luminous sources increases

In context of gamma signatures of UHECR :

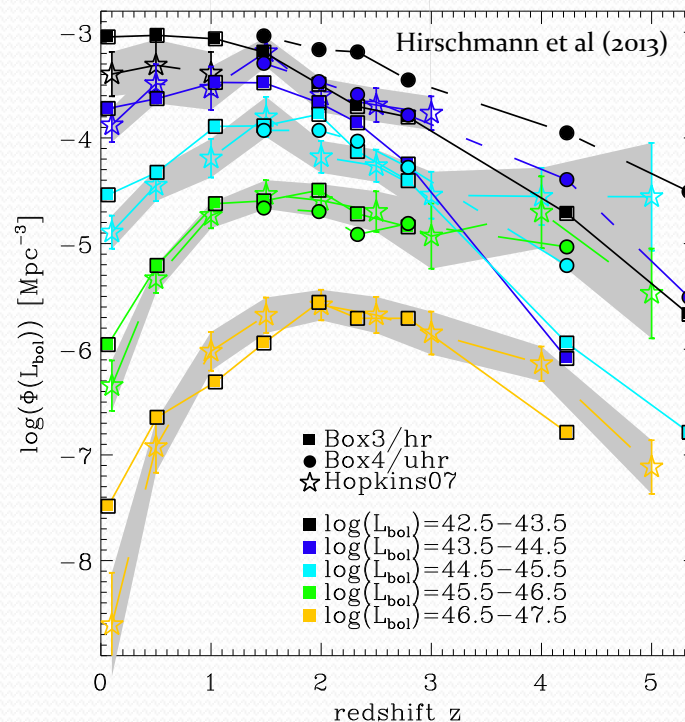
- Relaxed requirements for accelerator
- Large redshifts allow to discriminate CR – induced cascade from blazars



Shankar et al (2009)

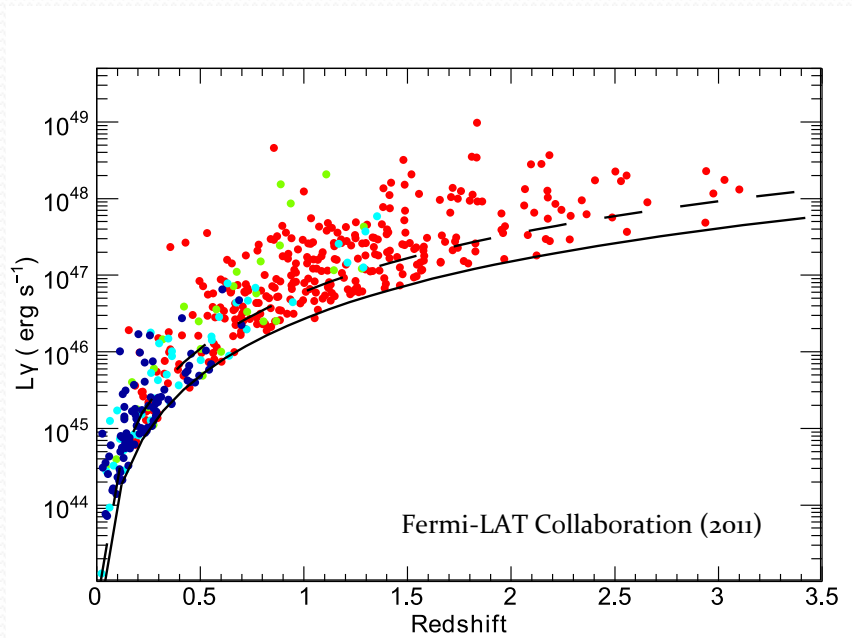
Distribution of AGNs with redshift

- The maximum of the distribution is for most luminous AGNs with $L=10^{45}$ - 10^{48} erg/s is at redshift $z=2$ -3.

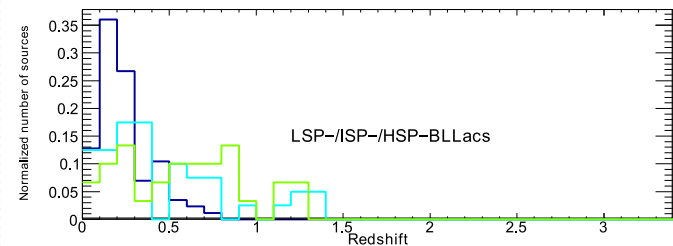
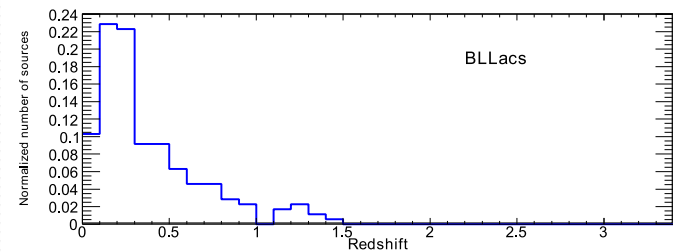
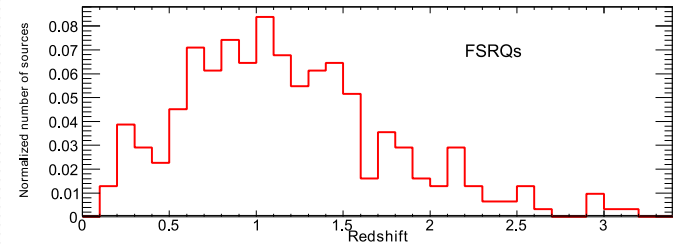


Source type distribution

- The Second Catalog of Active Galactic Nuclei



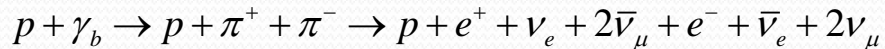
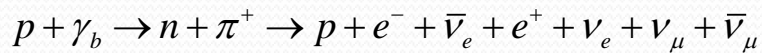
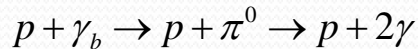
- More distant and luminous are FSRQs
- Small part of BLLacs at $z \sim 1$



Gamma-ray traces of UHECR

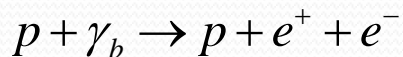
- Intergalactic medium: CMB+EBL

- Photomeson production (mostly CMB)



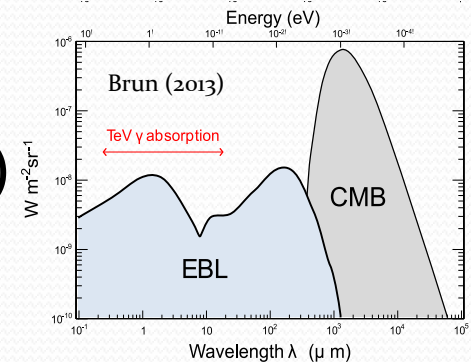
- Final products $E_\gamma \sim 10^{19} E_{p,20} eV$ $E_{e^\pm} \sim E_\nu \sim 5 \cdot 10^{18} E_{p,20} eV$

- Bethe-Heitler pair production (CMB+EBL)



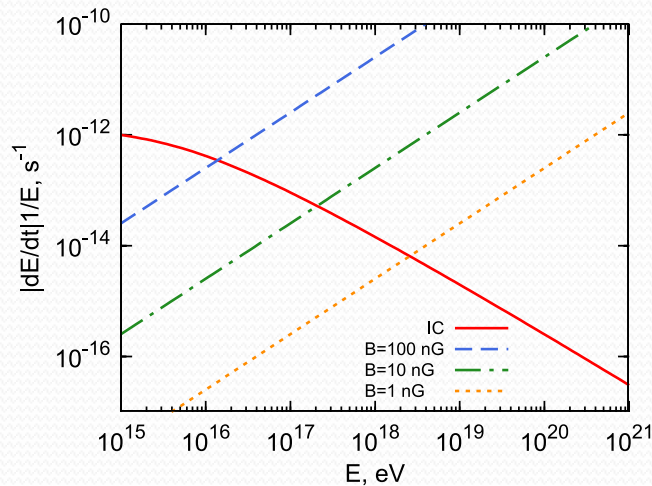
- Final products $E_{e^\pm} \sim 10^{15} E_{p,18} eV$

- Pair production $\gamma + \gamma_b \rightarrow e^+ + e^-$

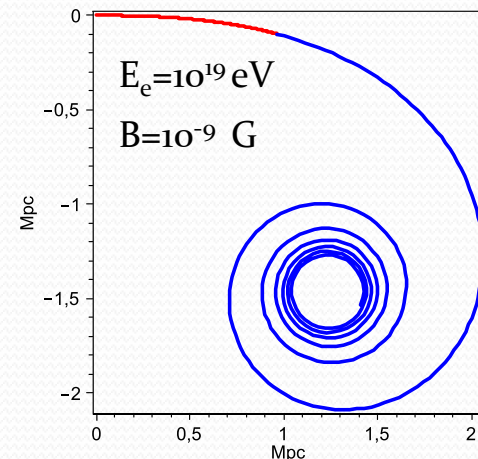


Two scenarios 1

- Synchrotron losses $>$ IC losses (10^{-10} G $<$ B $<$ 10^{-7} G)
 - Energy of electrons is effectively converted to synchrotron gamma rays of GeV and TeV energy band
 - Almost all energy is radiated on **the initial part of path**



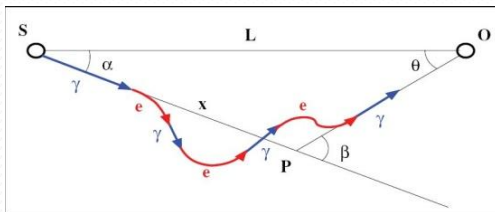
Energy loss rates of electrons



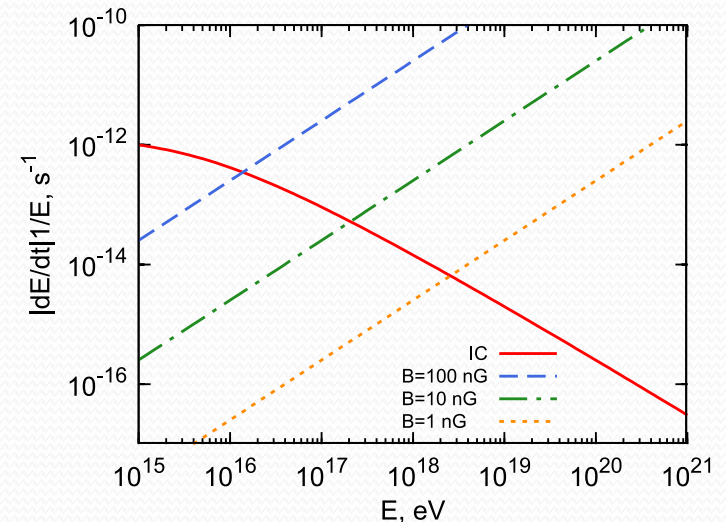
Illustrative: movement of electron in the constant magnetic field

Two scenarios 2

- IC losses > Synchrotron losses
 - Electromagnetic cascade: $\gamma + \gamma_b \rightarrow e^+ + e^- \rightarrow e^\pm + \gamma_b \rightarrow e^\pm + \gamma$
 - Magnetic field $10^{-14} \text{ G} < B < 10^{-10} \text{ G}$:
 - Strong deflection
 - Hardly detectable gamma ray halo
 - Magnetic field $B < 10^{-14} \text{ G}$:
 - Collimated beam of gamma rays



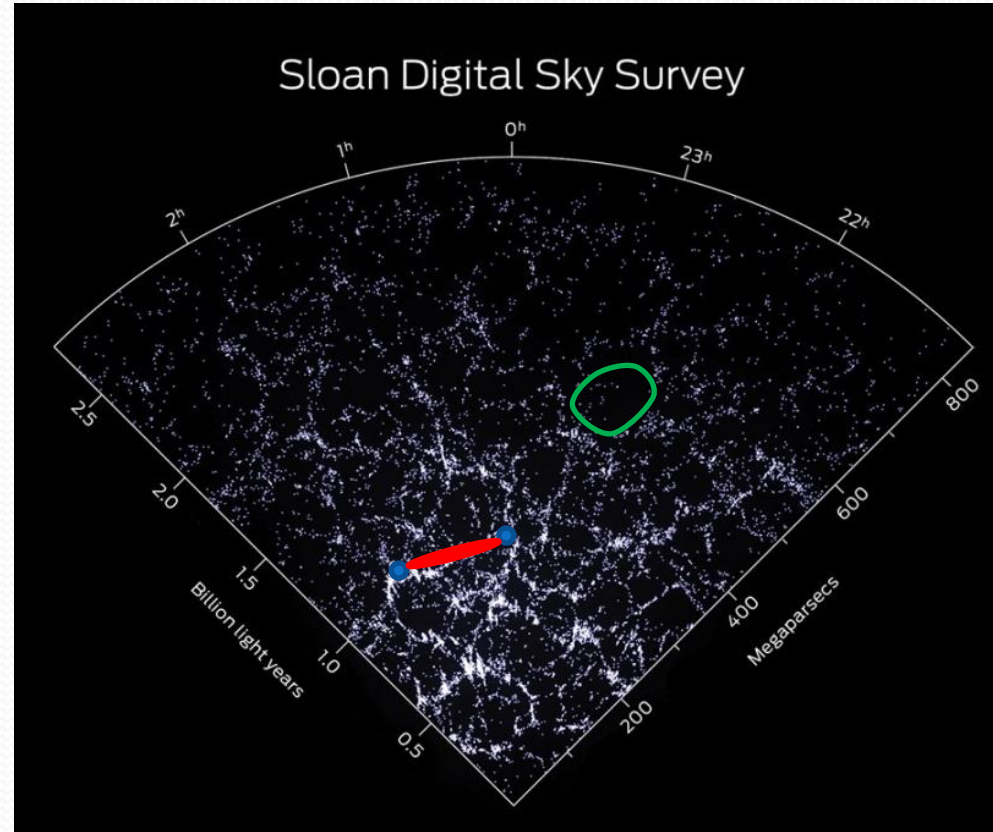
Dolag et al. (2009)



Energy loss rates of electrons

Intergalactic magnetic fields

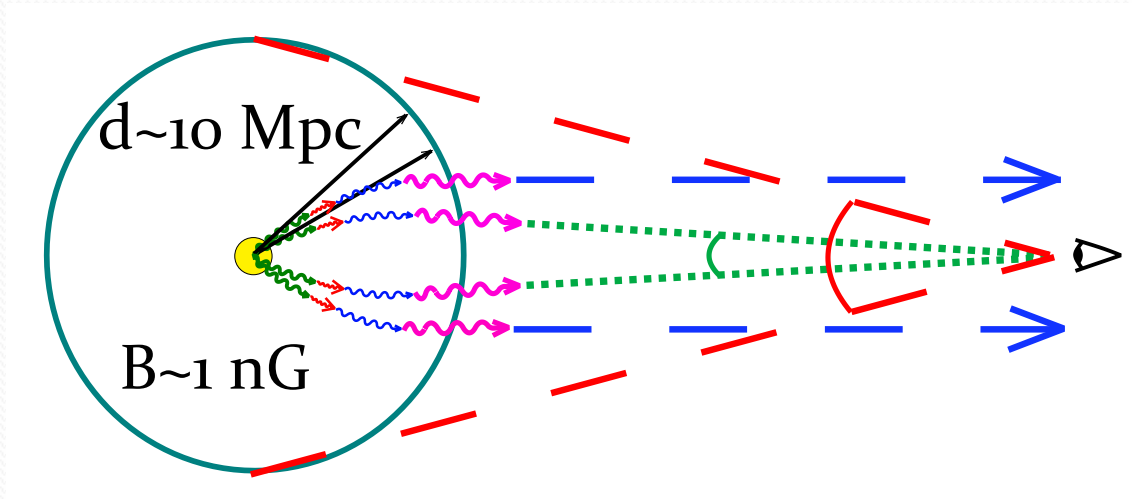
- **Clusters of galaxies**
 - 2 - 10 Mpc
 - $B=10^{-6}$ G
 - Faraday rotation
- **Filaments**
 - 50-80 Mpc
 - $B=10^{-9} - 10^{-8}$ G
 - Low frequency synchrotron radiation
- **Voids**
 - 10-150 Mpc
 - $B=10^{-17} - 10^{-15}$ G
 - Electromagnetic cascades from blazars



Scenario 1

Synchrotron radiation from the source of UHECR located at environment with strong magnetic field $10^{-10} \text{ G} < B < 10^{-7} \text{ G}$ (filaments)

Overview



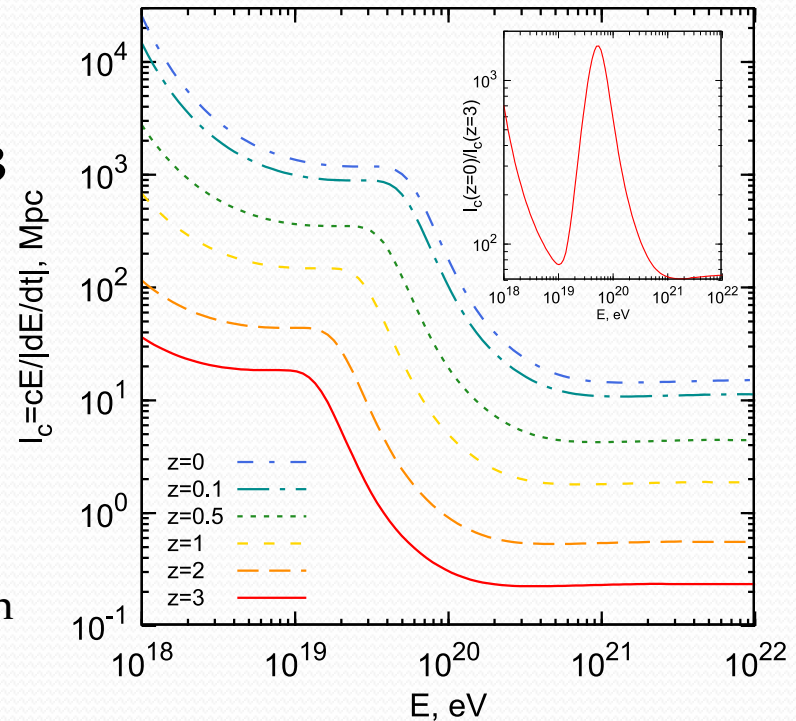
- Small deflection angles of protons and electrons result in collimated beam of synchrotron gamma rays
- Apparent angular size of gamma-ray source is smaller than angular size of the radiation region
- Non-variable source

Interaction of protons at large redshifts

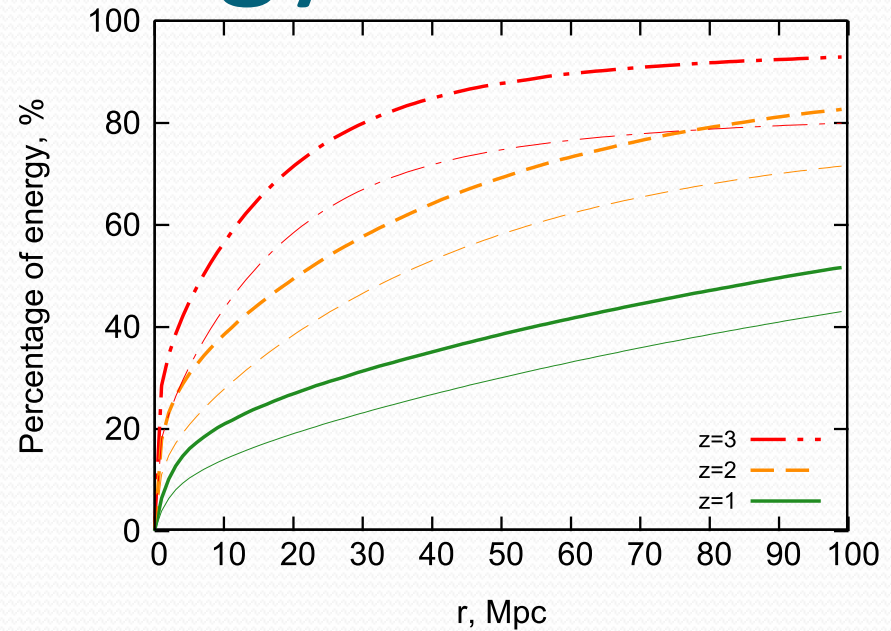
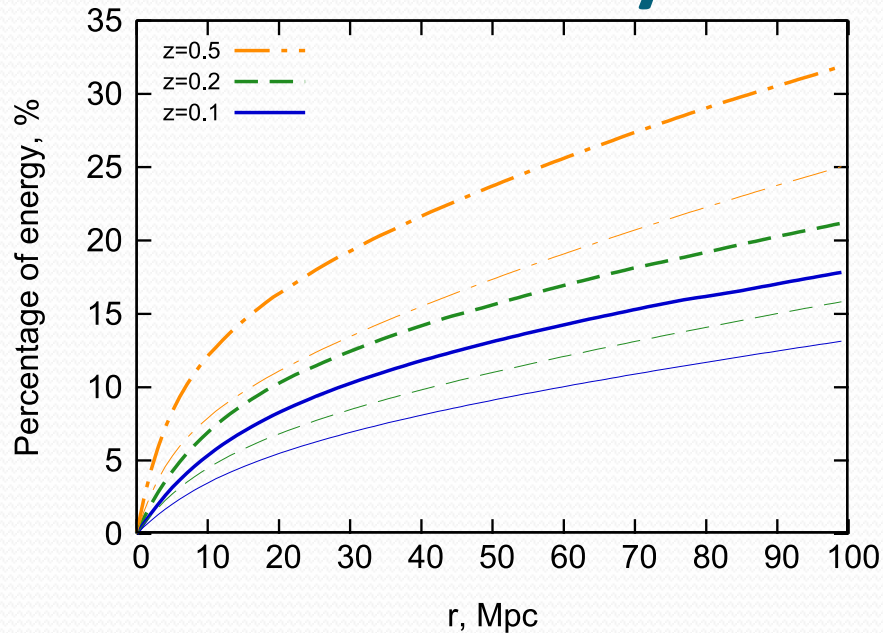
- Energy loss of UHE protons at different redshifts changes according

$$\beta(E, z) = (1+z)^3 \beta_0((1+z)E)$$

- Faster energy losses $\sim (1+z)^3$
 - Interaction lengths are shorter:
from 15 Mpc at $z=0$ to 0.2 Mpc at $z=3$
- Interaction with more energetic photons $\sim 1+z$
 - Less energetic protons produces energetic electrons via photo-hadronic interactions
 - Relaxed requirements for acceleration
 - Increase number of UHECR sources seen in gamma rays

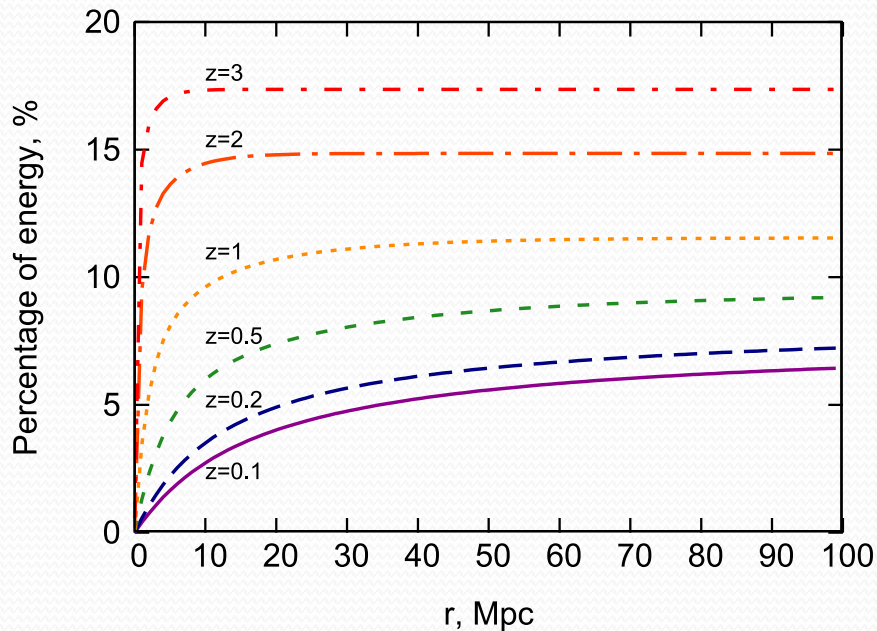


Efficiency of energy transfer

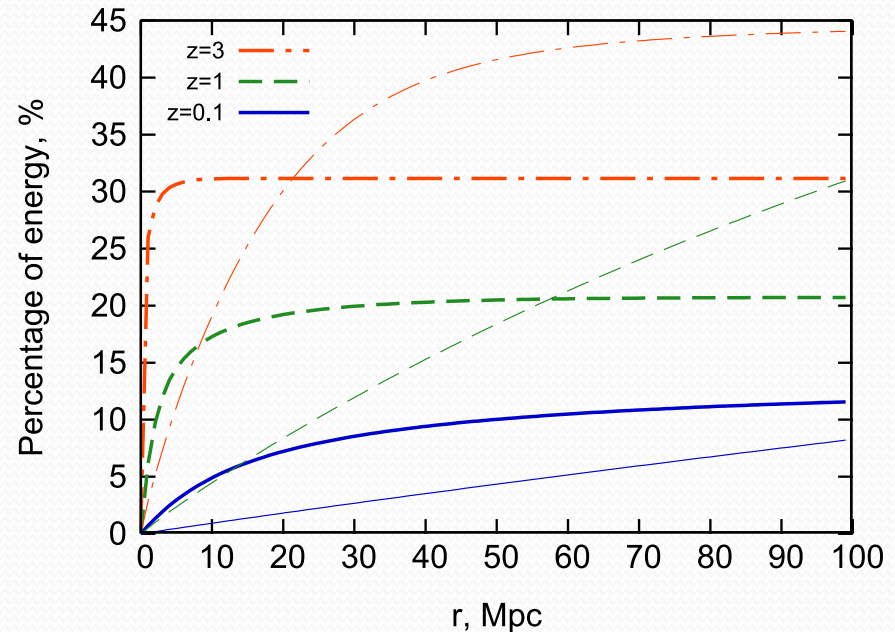


- Energy losses of protons with $E > 10^{18}$ eV increases from 19% to 93%
- Energy transferred to the electrons increases from 13% to 80%

Efficiency of energy transfer



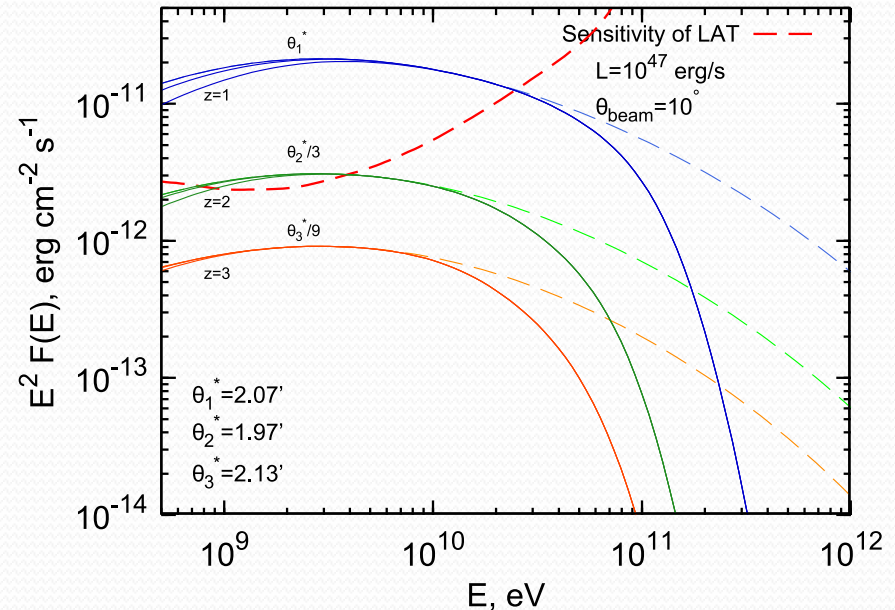
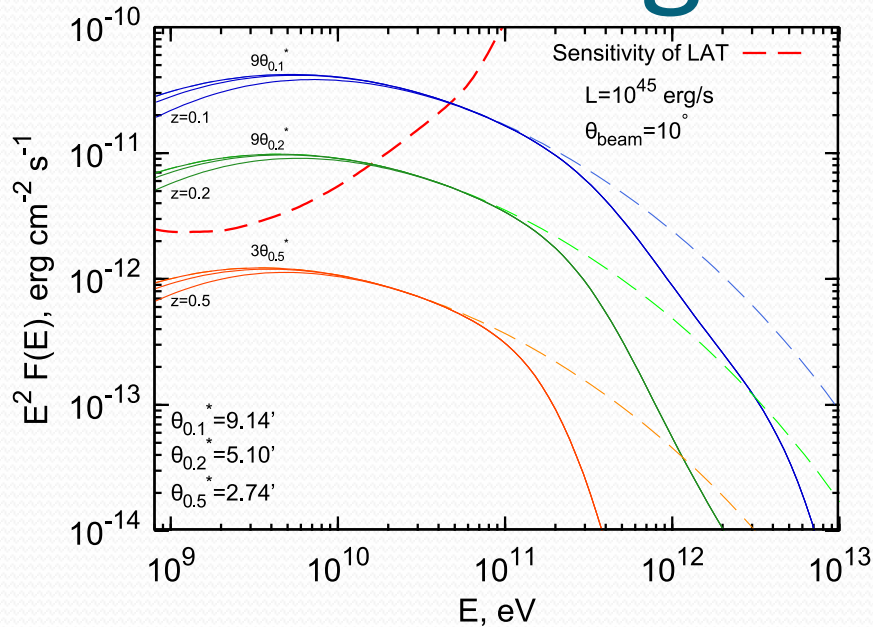
For protons with $E > 10^{18}$ eV



For protons with $E > 10^{19}$ eV

- All the possible energy is transferred at first 5 Mpc at $z=3$
- Considerable part of the energy goes to electrons produced via pair production

SED of gamma-ray signal



- Flux of gamma rays is detectable by Fermi LAT for isotropic luminosity $L=10^{45} - 10^{47}$ erg/s
- For sources at $z > 0.2$ the gamma ray source is point like for Fermi LAT
- Flux is suppressed at TeV energies
- With increase of strength of magnetic field spectra shifts to higher energy proportionally to increase of magnetic field strength
 - At large magnetic fields energetic part of the spectra will be absorbed due to interaction with EBL

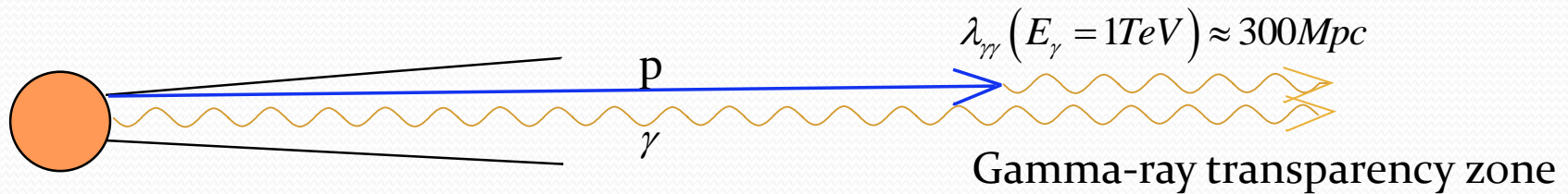
Scenario 2

Gamma rays from electromagnetic cascade induced by cosmic rays
in the environment with weak magnetic field $10^{-17} \text{ G} < B < 10^{-15} \text{ G}$ (voids)

Distant TeV blazars

- The flux of blazars with redshift $z > 0.1$ at TeV energies should be significantly suppressed due to gamma-ray absorption in EBL
- Recovered intrinsic spectrum $F_{\text{obs}}(E)\exp(\tau(E,z))$ some of them (1ES 0229+200 ($z=0.139$), 1ES 0347+121 ($z=0.185$), 1ES 1101+232 ($z=0.186$)) seems harder than it is expected in standard models with photon index $\Gamma > 1.5$
- Explanations
 - Intrinsic properties of the radiation
 - Gamma-ray induced cascade
 - Cosmic-ray induced cascade
- Cosmological distant blazars observed at TeV energies the only feasible explanation is cosmic-ray induced cascade

Blazars as CR sources



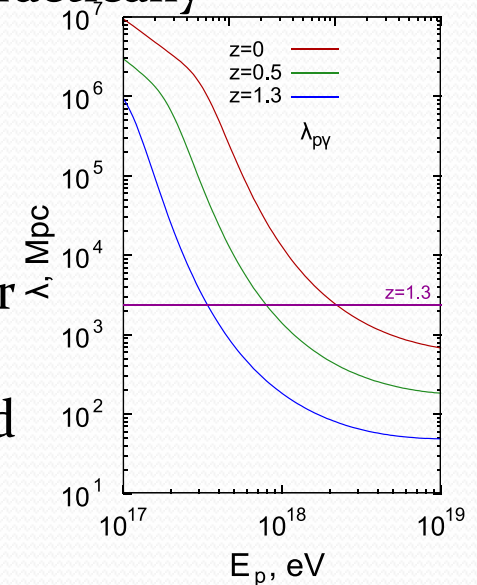
- Protons with energies $E=10^{17}\text{-}10^{19}\text{ eV}$ has large interaction length relative to pair production process
- In the magnetic field $B<10^{-15}\text{ G}$ protons propagates practically rectilinearly

$$\theta_p \approx 0.05\text{arcmin} \left(\frac{10^{18}\text{ eV}}{E_p} \right) \left(\frac{B}{10^{-15}\text{ G}} \right) \left(\frac{L}{\text{Mpc}} \frac{d}{\text{Gpc}} \right)^{1/2}$$

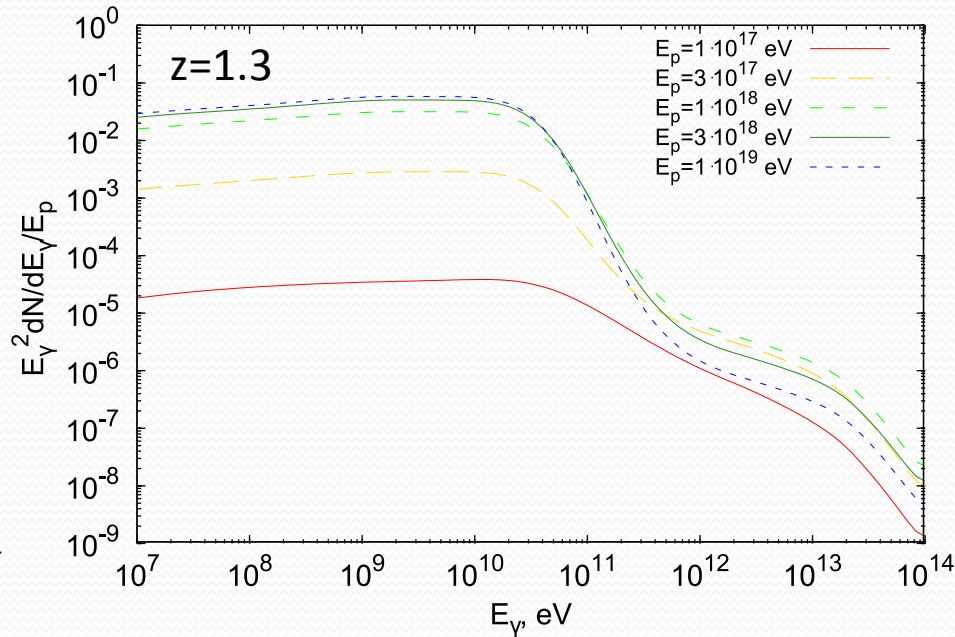
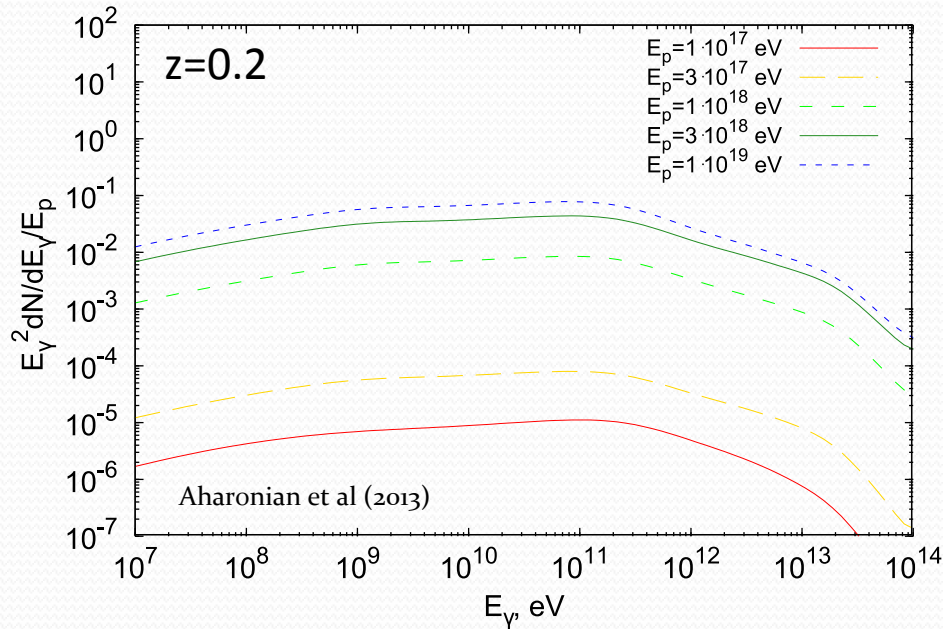
- Mean free path without encounter with galaxy cluster

$$L_{\min} \sim 1 / (\pi R^2 n) \sim (1-5) \times 10^3 \text{ Mpc}$$

- Encounter of the structure with strong magnetic field significantly suppresses the flux of gamma rays

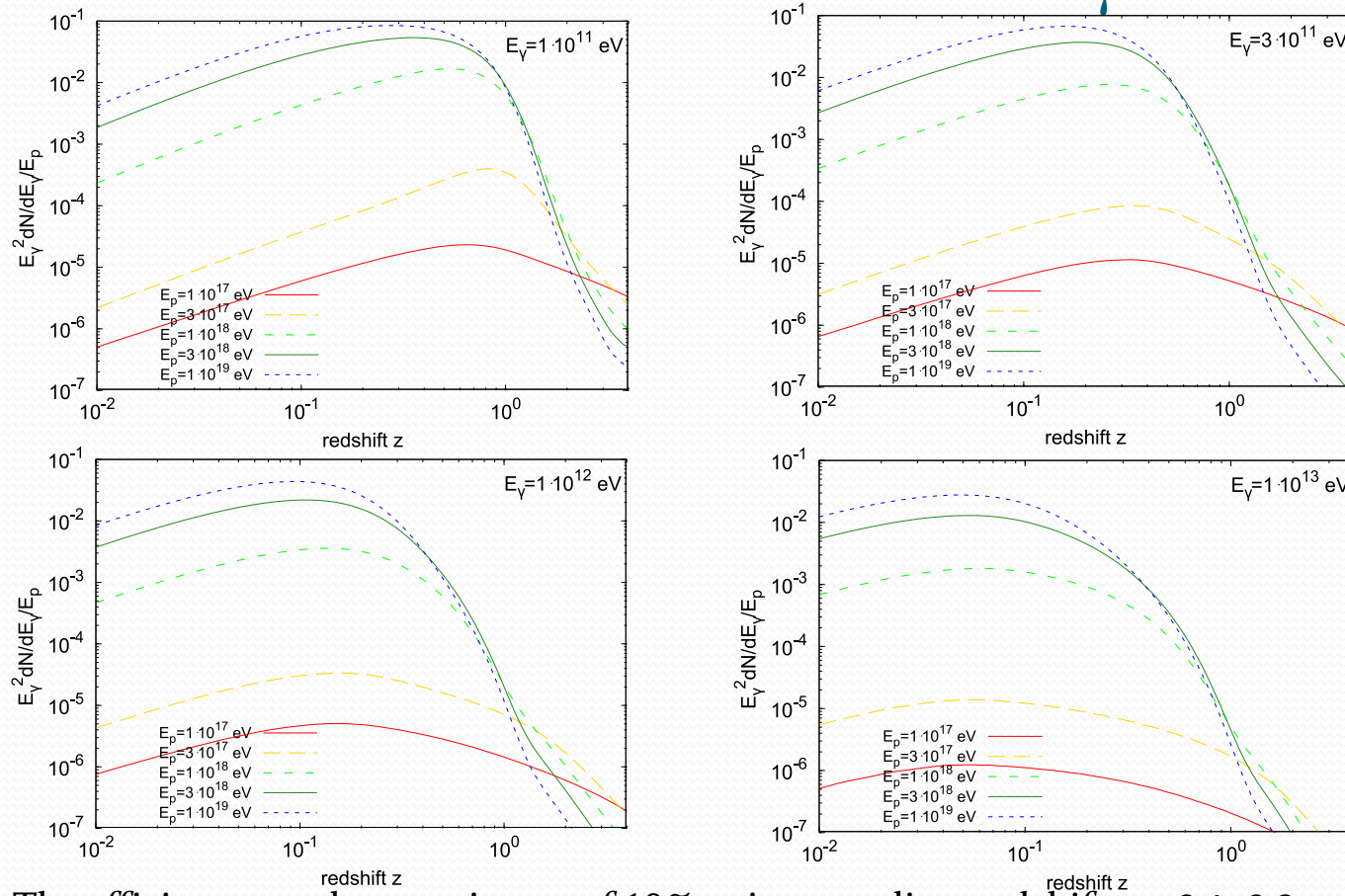


Cascade spectra at small and large redshifts



- For small redshifts spectral shape of secondary photons does not depend on the initial energy of protons with maximum at $E=10^{11}$ eV
 - For a nearby source the spectral shape of secondary photons is independent of the details of the proton energy spectrum
- For large redshifts spectral shape of secondary photons significantly depends on initial energy of protons with hardening at $E>1$ TeV
- For large redshifts mostly the low-energy protons with $E \sim 10^{17}$ eV enter the gamma-ray transparency zone

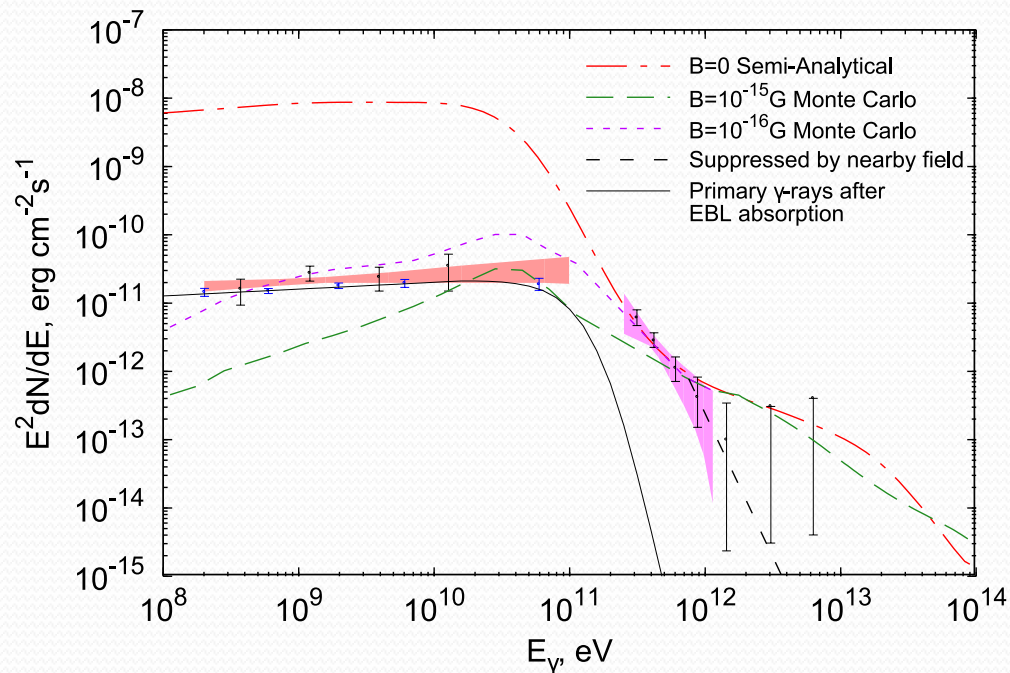
Efficiency of energy transfer to the flux of gamma rays with energy E_γ



- The efficiency reaches maximum of 10% at intermediate redshifts $z = 0.1-0.3$
- The efficiency falls down to 10^{-4} - 10^{-5} at redshift $z \sim 1$
- The most efficient transfer of energy for protons with initial energy $E = 10^{18}$ eV

Case of PKS 0447-439 at $z=1.3$

- In spite on the low efficiency protons still can explain TeV radiation from cosmologically distant blazars if $B < 10^{-15}$ G
- The suppression of flux at GeV energies due to deflection of electrons in the electromagnetic cascade
- To fit the Fermi data the magnetic field should be $B \sim 10^{-16}$ G
- If magnetic is strong ($B > 10^{-14}$ G) in the vicinity of the observer (~ 100 Mpc), the flux is suppressed at TeV energies



Implications

- Deflection of protons $\theta_p \approx 0.05 \text{ arcmin} \left(\frac{10^{18} \text{ eV}}{E_p} \right) \left(\frac{B}{10^{-15} \text{ G}} \right) \left(\frac{L}{\text{Mpc}} \frac{d}{\text{Gpc}} \right)^{1/2}$
- Deflection in the cascade $\theta_{cas} \approx 3.8 \text{ arcmin} \left(\frac{10^{12} \text{ eV}}{E_\gamma} \right) \left(\frac{B}{10^{-15} \text{ G}} \right)$
 - The main contribution to deflection is from cascade
 - Energy-dependent angular broadening can provide a direct measurements of IGMF in a given direction
- Time delay of protons $\Delta\tau_p \approx 1.5 \cdot 10^6 \text{ s} \left(\frac{E_p}{10^{18} \text{ eV}} \right)^{-2} \left(\frac{B}{10^{-15} \text{ G}} \right)^2 \left(\frac{L}{1 \text{ Mpc}} \right) \left(\frac{d}{1 \text{ Gpc}} \right)^2$
- Time delay in the cascade $\Delta\tau_\gamma \approx 1.3 \cdot 10^6 \text{ s} \left(\frac{E_\gamma}{10^{12} \text{ eV}} \right)^{-5/2} \left(\frac{B}{10^{-15} \text{ G}} \right)^2$
 - For $B \sim 10^{-15} \text{ G}$ any variability of the initial signal on the scales of the order 1 month and shorter is smeared out
 - In framework of this model any observed variability shorter than month means $B < 10^{-15} \text{ G}$
- Isotropic luminosity $L = 10^{50} \text{ erg/s}$ is needed
- Luminosity with beaming $L = 10^{47} \text{ erg/s}$ which is of the order of Eddington luminosity of black hole with mass $M \sim 10^9 M_{\text{sol}}$

Conclusion

- As AGNs is more active at the epoch of $z=2..3$ there are more accelerators of UHECR
- The discovery of UHECR sources is possible through their gamma rays signatures
- Two significantly different cases should be distinguished:
- Synchrotron gamma rays
 - Strong nG level magnetic field, short interaction lengths
 - At large redshifts relaxed condition for accelerators
- Gamma rays produced through cascade
 - Weak fG level magnetic field, large interaction length
 - At large redshifts only cosmic rays production in blazars can explain TeV gamma rays